

# Scaling and Finite-Size Scaling analysis of critical behavior in lattice QCD

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J. Braun and B. Klein, Phys. Rev. D **77** (2008) 096008;  
PoS **LAT2007**:198 (2007) and in preparation (2008).

# Overview

Introduction

Scaling in Infinite volume

Finite-Size Scaling

Strategy for comparison

Conclusions

# Phase transitions in QCD

## QCD phase transitions

- ▶ de-confinement phase transition
- ▶ chiral phase transition

# Phase transitions in QCD

## QCD phase transitions

- ▶ de-confinement phase transition
- ▶ chiral phase transition

## Lattice Gauge Theory

- ▶ fully non-perturbative method
- ▶ finite simulation volume
- ▶ explicit symmetry breaking through quark masses
- ▶ phase transition order? → (finite-size) scaling analysis



# Phase transitions in Finite Volume

- ▶ no actual phase transition in finite volume
- ▶ thermodynamic potential completely regular
- ▶ order parameter vanishes strictly . . .
- ▶ . . . unless explicit symmetry breaking is present
- ▶ critical behavior difficult to identify

finite-size scaling is an essential tool for the analysis  
of critical behavior in finite-volume simulations

- $N_f = 2$ : second order for  $m_q = 0$ , crossover for  $m_q \neq 0$ ,  $O(4)$

R. D. Pisarski and F. Wilczek, Phys. Rev. D **29** (1984) 338.

- first order phase transition? (confinement dominates?)  
(staggered fermions)

M. D'Elia, A. Di Giacomo and C. Pica, Phys. Rev. D **72** (2005) 114510 [arXiv:hep-lat/0503030];

G. Cossu, M. D'Elia, A. Di Giacomo, and C. Pica (2007), arXiv:0706.4470 [hep-lat].

- consistent with  $O(2)$  for  $\chi$ QCD (staggered) (but no scaling)

J. B. Kogut and D. K. Sinclair, Phys. Rev. D **73** (2006) 074512 [arXiv:hep-lat/0603021].

- decide by analyzing scaling behavior

## A scaling analysis requires *critical exponents* and *scaling functions*

- ▶ Scaling functions obtained mostly from  $O(N)$  lattice simulations and perturbative RG

D. Toussaint, Phys. Rev. D55 (1997) 362.

J. Engels, S. Holtmann, T. Mendes, and T. Schulze, Phys. Lett. **B514** (2001) 299.

E. Brézin, D. J. Wallace, and K. Wilson, Phys. Rev. B7, 232 (1973).

F. Parisen Toldin, A. Pelissetto and E. Vicari, JHEP **0307** (2003) 029.

- ▶  $O(N)$  critical exponents from FRG calculations

N. Tetradis and C. Wetterich, Nucl. Phys. **B422** (1994) 541.

O. Bohr, B.J. Schaefer, and J. Wambach, Int. J. Mod. Phys. **A16** (2001) 3823.

D. F. Litim and J. M. Pawłowski, Phys. Lett. B **516** (2001) 197.

- ▶ few results with explicit symmetry breaking
- ▶ no results on finite-size scaling

# Functional RG for the $O(N)$ -model

Effective action at scale  $k$  ( $\Lambda \geq k \geq 0$ ) in LPA ( $\eta = 0$ )

$$\Gamma_k[\phi] = \int d^d x \frac{1}{2} (\partial_\mu \phi)^2 + \\ + a_1(k) (\phi^2 - \sigma_0^2(k)) + a_2(k) (\phi^2 - \sigma_0^2(k))^2 + \dots - H(\sigma - \sigma_0(k))$$
$$\phi = (\sigma, \vec{\pi}) \quad \phi^2 = \sigma^2 + \vec{\pi}^2 \quad O(N)\text{-symmetric}$$

scale-dependent couplings

$$\sigma_0(k), a_n(k), \quad n = 1, \dots, n_{\max} \quad 2a_1(k)\sigma_0(k) = H \equiv \text{const.}$$



## Inputs:

- ▶ choice of initial scale  $\Lambda$  sets all scales
- ▶ couplings  $\sigma_0(\Lambda)$  and  $a_2(\Lambda)$  at initial scale  $\Lambda$
- ▶  $d = 3$ : no field-theoretical temperature
- ▶  $\sigma_0(\Lambda) - \sigma_0^{\text{crit}}(\Lambda) \sim T - T_c$ :  
     $\sigma_0(\Lambda) > \sigma_0^{\text{crit}}(\Lambda) \Rightarrow$  broken phase  
     $\sigma_0(\Lambda) < \sigma_0^{\text{crit}}(\Lambda) \Rightarrow$  symmetric phase

application to **finite volume**:

$$\int \frac{d^d p}{(2\pi)^d} f(p^2) \longrightarrow \frac{1}{L^d} \sum_{n_1, \dots, n_d} f\left(\left(\frac{2\pi}{L}\right)^2 (n_1^2 + \dots + n_d^2)\right)$$

- ▶ local expansion around the minimum
- ▶ lowest order in local potential approximation:  $\eta = 0$
- ▶ potential with explicit symmetry breaking
- ▶ ERG with optimized cutoff D. F. Litim, Phys. Lett. B486 (2000) 92.  
     $\Leftrightarrow$  proper-time RG (infinite volume) S. B. Liao, Phys. Rev. D53 (1996) 2020.  
    D. F. Litim and J. M. Pawłowski, Phys. Lett. B 516 (2001) 197.

# Scaling for the singular free energy

there is no length scale at a critical point:  $\xi \rightarrow \infty$

- ▷ scale invariance of free energy density at critical point
- ▷ use behavior close to critical point

$$f_s(t, h) = \ell^{-d} f_s(\ell^{y_t} t, \ell^{y_h} h)$$

$$t = (T - T_c)/T_0, \quad h = H/H_0, \quad y_t = \frac{1}{\nu}, \quad y_h = \frac{\nu}{\beta\delta}$$

idea: keep one of the arguments fixed

- ▷ becomes function of a single scaling variable

# Scaling function in Fisher parametrization

$$M \sim \frac{\partial}{\partial H} f(t, h)$$

$$t = (T - T_c)/T_0, \quad h = H/H_0, \quad z = \frac{t}{h^{1/(\beta\delta)}}, \quad \xi(t) \sim t^{-\nu}$$

## Scaling function for the order parameter $M$

$$\left. \begin{array}{l} M(t, h=0) = (-t)^\beta \\ M(t=0, h) = h^{1/\delta} \end{array} \right\} \rightarrow M(t, h) = h^{1/\delta} f(z), \quad f(z) \xrightarrow{z \rightarrow -\infty} (-z)^\beta$$

critical exponent  $\beta$  enters into asymptotic behavior

# Scaling function in Fisher parameterization

## Scaling function for the susceptibility $\chi$

$$\chi = \frac{\partial M}{\partial H} \Rightarrow \chi = \frac{1}{H_0} h^{1/\delta-1} \frac{1}{\delta} \left[ f(z) - \frac{z}{\beta} f'(z) \right] = \frac{1}{H_0} h^{1/\delta-1} f_\chi(z)$$

▷ consequence of scaling:

$f_\chi(z)$  completely specified by  $f(z)$

# Critical exponents for the O(4) model in d=3

	$\nu$	$\beta$	$\eta$	$\delta$
lattice <sup>1</sup>	0.7423	0.380	0.024*	4.86
RG <sup>2</sup>	0.8043	0.4022*	—	5.00*
RG	0.8053(60)	0.4030(30)	—	4.973(30)

<sup>1</sup> J. Engels, S. Holtmann, T. Mendes and T. Schulze, Phys. Lett. B **514** (2001) 299.

<sup>2</sup> D. F. Litim and J. M. Pawłowski, Phys. Lett. B **516** (2001) 197.

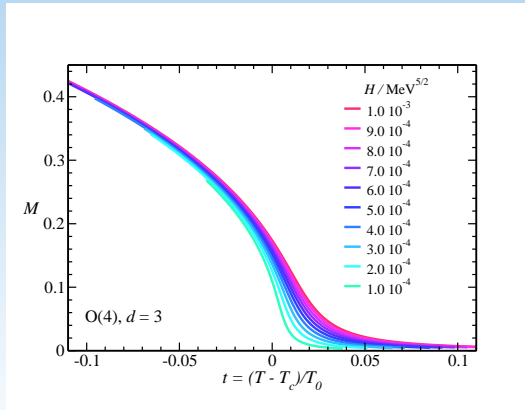
If we already know the critical exponents, why do we want the scaling functions as well?

$$\beta\delta/\nu = (\gamma + \beta)/\nu = (2 - \eta) + \frac{1}{2}(d - 2 + \eta) \approx \frac{5}{2}$$

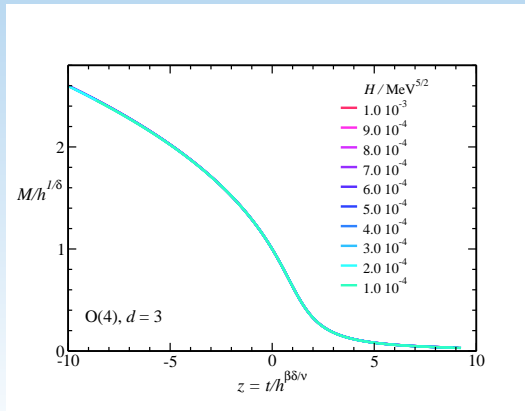
$$\beta/\nu = \frac{1}{2}(d - 2 + \eta) \approx \frac{1}{2}$$

# Order parameter $M$ as a function of $t$

for small values of  $H$

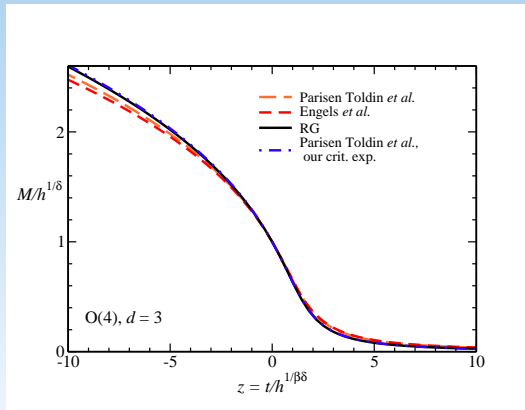


# Rescaled order parameter $M/h^{1/\delta}$ as a function of $z$ for small values of $H$





# Comparison to lattice spin-model $O(4)$ result

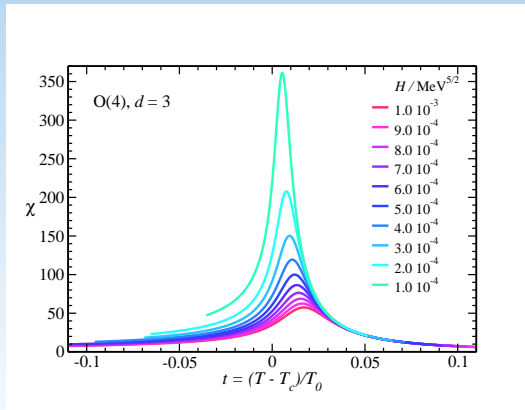


J. Engels, S. Holtmann, T. Mendes, and T. Schulze, Phys. Lett. **B514** (2001) 299

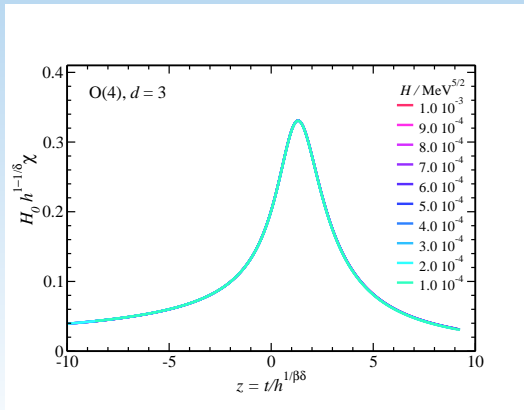
F. Parisen Toldin, A. Pelissetto and E. Vicari, JHEP **0307** (2003) 029.

# Susceptibility $\chi$ as a function of $t$

for small values of  $H$  (note trajectory  $t_p/h^{1/(\beta\delta)} = z_p$ )

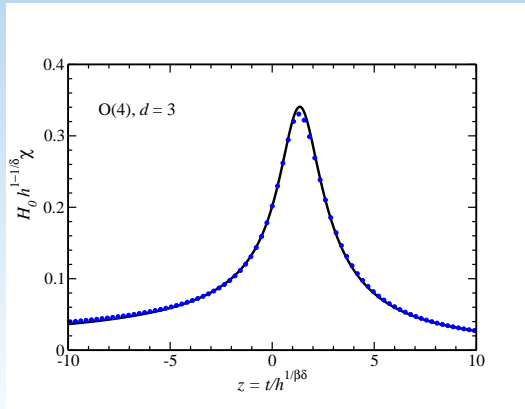


# Rescaled Susceptibility $H_0 h^{1-1/\delta} \chi$ as a function of $z$ for small values of $H$



## Test of scaling function

Is it true that  $f_\chi(z) = \frac{1}{\delta} \left[ f(z) - \frac{z}{\beta} f'(z) \right]$ ? Yes!



# Finite-Size Scaling

- ▶ Universal behavior requires divergence of the correlation length  $\xi$
- ▶ Finite volume size  $L$  cuts off the critical fluctuations
- ▶ Universal scaling behavior is therefore affected if correlation length  $\xi \sim L$ , depends on ratio  $\xi/L$

Finite-Size Scaling hypothesis (Fisher): The ratio of thermodynamic quantities ( $M, \chi, \dots$ ) in the finite-size system and the infinite-size system is a function of *only* the ratio  $\xi/L$ :

$$\frac{M_L(t)}{M_\infty(t)} = \mathcal{F}\left(\frac{L}{\xi(t)}\right), \quad \xi(t) \sim t^{-\nu}, \quad M(t, h) = h^{1/\delta} f(z)$$

# Finite-Size Scaling Functions

Idea for obtaining the universal Finite-Size Scaling functions:

- ▶ keep  $L/\xi = \text{const.} \rightarrow \text{vary } t \sim L^{1/\nu}$
- ▶ keep  $z = t/h^{1/(\beta\delta)} = \text{const.} \rightarrow \text{vary } h \sim L^{-\beta\delta/\nu}$

$$M(t, h) = h^{1/\delta} f(z) \rightarrow L^{-\beta/\nu} (h L^{\beta\delta/\nu})^{1/\delta} f(z)$$

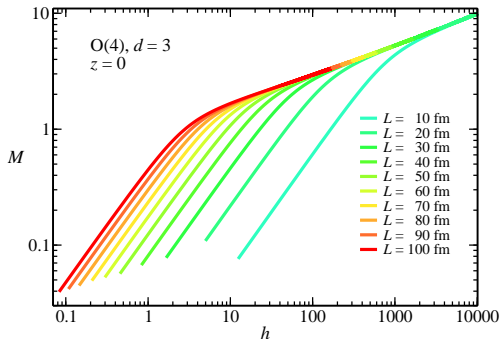
Finite-Size Scaling Functions depend only on  $h L^{\beta\delta/\nu}$   
(for any given value of  $z$ ):

$$L^{\beta/\nu} M = Q_M(z, h L^{\beta\delta/\nu})$$

$$L^{\gamma/\nu} \chi = Q_\chi(z, h L^{\beta\delta/\nu})$$

# Finite-Size Scaling

Order parameter  $M(h)$  vs.  $h$  for  $L = 10 - 100$  fm

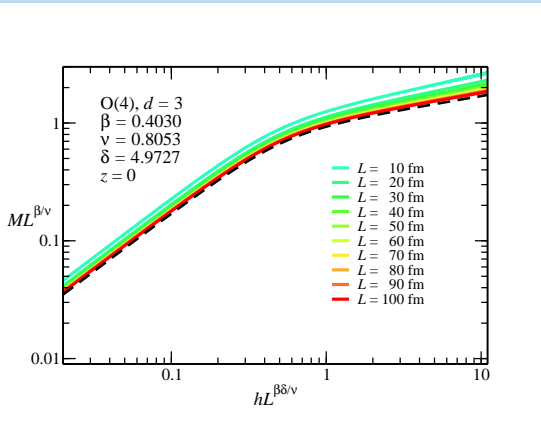


at the critical temperature ( $z = 0$ ).

- ▶  $\xi$  small for large  $h$
- ▶ deviations for  $\xi \sim L$
- ▶ asymptotic behavior given by  $1/\delta$

# Finite-Size Scaling

Finite-size scaled order parameter  $ML^{\beta/\nu}$  vs.  $hL^{\beta\delta/\nu}$

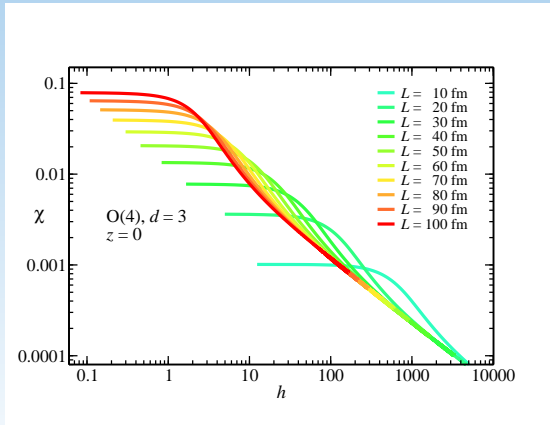


for  $L = 10 - 100$  fm.

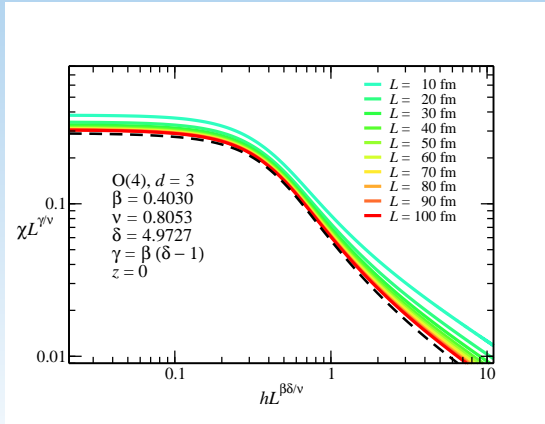
- ▶ scaling deviations for large fields  $h$
- ▶ controlled by sub-leading operator
- ▶ consistent with RG prediction for  $\omega$
- ▶ extrapolate to obtain scaling function



## Susceptibility $\chi(h)$ vs. $h$ for $L = 10 - 100$ fm

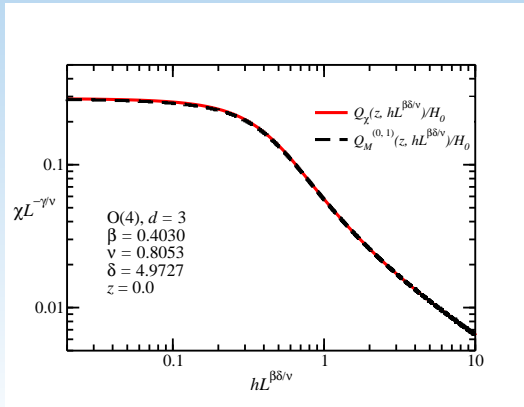


# Finite-size scaled susceptibility $\chi L^{\gamma/\nu}$ vs. $h L^{\beta\delta/\nu}$



# Test of the scaling function

Is it true that  $Q_\chi(z, hL^{\beta\delta/\nu}) = \frac{\partial}{\partial(hL^{\beta\delta/\nu})} Q_M(z, hL^{\beta\delta/\nu})$ ? Yes!



# Comparison to lattice QCD results

exploit *universality* near a critical point to analyze behavior

## Issues to address

- ▶ Comparison of universal scaling functions
- ▶ only IR quantities can be compared, different UV physics
- ▶ determine non-universal normalization factors *or*
- ▶ determine dimensionless ratios
- ▶ locate finite-size scaling regions
- ▶ estimate scaling corrections

# Scale normalizations

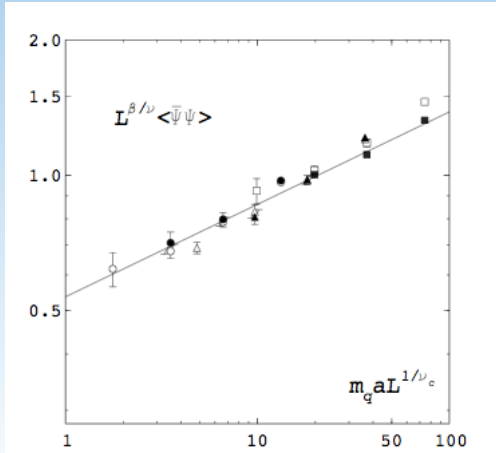
- ▶ for every coupling  $(t, h, \dots)$  there is a normalization scale  $(T_0, H_0, \dots)$
- ▶  $\bar{h} = hL^{\beta\delta/\nu}$  is not dimensionless  $\Rightarrow$  not universal
- ▶ need universal quantity: Length scale normalization required

best candidate: dimensionless ratio/product of IR quantities

$$\frac{L}{\xi(t, h, L)} = M_\sigma(t, h, L)L \geq M_\pi(t, h, L)L$$

$$\xi(t) = \frac{1}{C_0} t^{-\nu}, \quad M_\sigma(t) = C_0 t^\nu \Rightarrow \frac{L}{\xi(t, L \rightarrow \infty)} = \frac{C_0 L}{t^{-\nu}} = \ell t^\nu$$

# Asymptotic finite-size scaling on the lattice

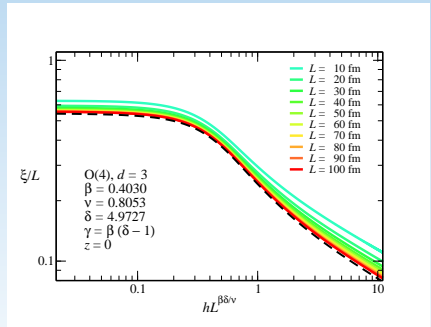
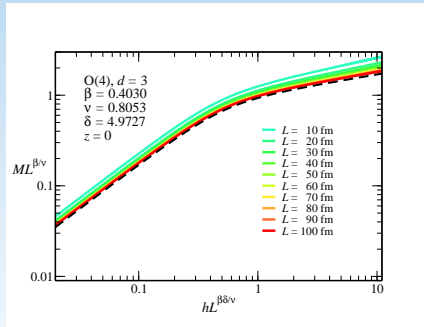


- ▶ asymptotic behavior for  $hL^{\beta\delta/\nu}$  large
- ▶ no large finite-size scaling effects in asymptotic region
- ▶ lattice- $N_s$  is dimensionless, but  $N_s = L/a(O(N)) \neq N_s = L/a(QCD)$ : need to fix  $\xi/L$

J. Engels, S. Holtmann, T. Mendes, and T. Schulze, Phys. Lett. **B514** (2001) 299.

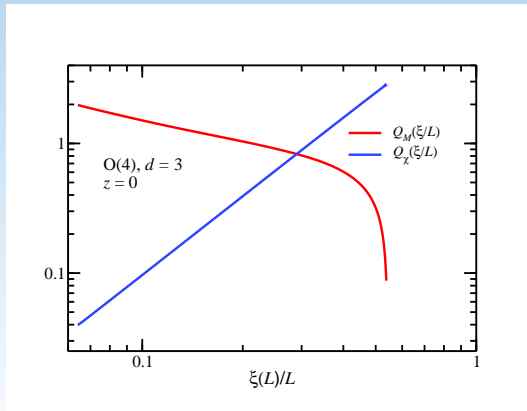
# Finite-Size Scaling regions

- pion mass provides bound for  $\xi/L$ :  $m_\pi \leq \frac{1}{\xi} \Leftrightarrow m_\pi L \leq \frac{L}{\xi}$
- $m_\pi L$  must not be too large to see sizable effects



# Observables as a function of $\xi/L$

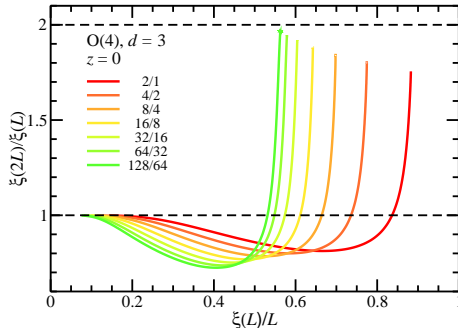
scaling functions plotted against the dimensionless ratio  $\xi(L)/L$





# Ratios of results from systems with fixed size ratio

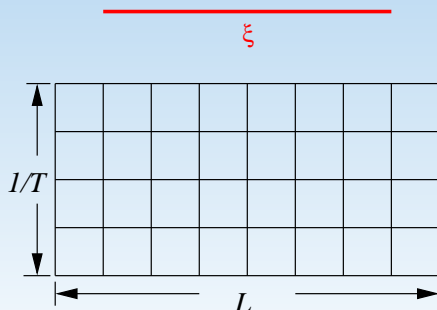
$\xi(2L)/\xi(L)$  as a function of  $\xi(L)/L$



- ▶ no normalization problems
- ▶ but large corrections
- ▶ correlation length bounded  $\xi(L) \leq k_0 L$
- ▶ ratio approaches volume ratio for  $L \rightarrow \infty$

# Dimensional reduction

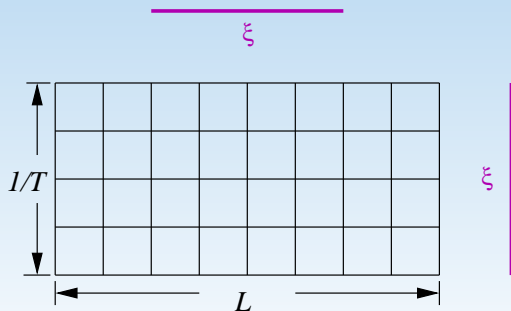
When can we expect scaling behavior in a  $d = 3$  class?



- ▶ for  $d = 3$  behavior need  $\xi \gg 1/T$
- ▶ for infinite-volume scaling behavior need  $\xi/L \ll 1$
- ▶ for finite-size scaling need  $\xi/L \approx 1$
- ▶  $m_\pi L \leq L/\xi$  is bound

# Dimensional reduction

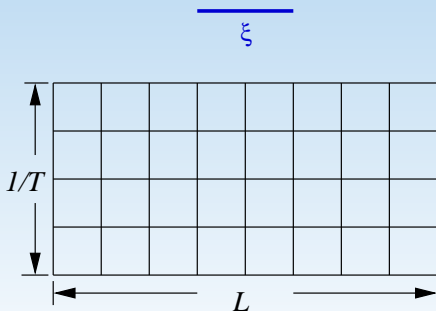
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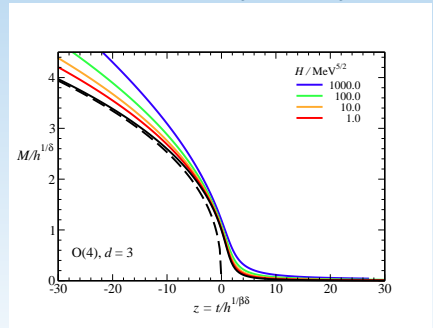
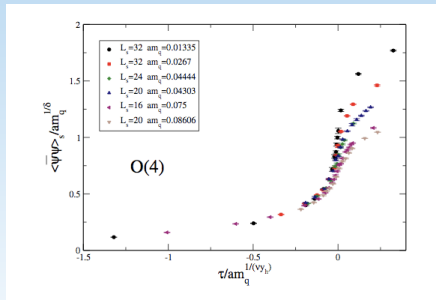
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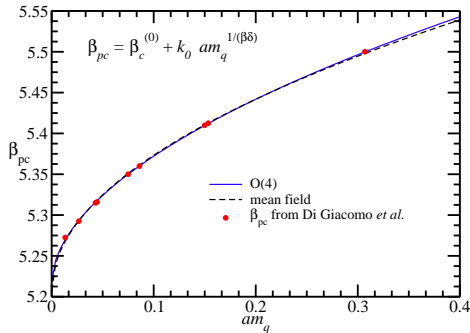
# Scaling of the order parameter on the lattice

- ▶ rescaled order parameter (subtracted) from DiGiacomo et al.
- ▶ scaling corrections largest in phase with broken symmetry



M. D'Elia, A. Di Giacomo and C. Pica, Phys. Rev. D **72** (2005) 114510 [arXiv:hep-lat/0503030]

# pseudocritical coupling on the lattice



what is the relevant  $T$ ?

$$T = \frac{1}{N_t a(\beta, am_q)}$$

$$\Rightarrow \frac{T_c - T}{T_c} = 1 - \frac{a(\beta_c^{(0)}, 0)}{a(\beta, am_q)}$$

what is the relevant  $h$ ?

$$h = am_q N_t$$

( $h = \beta H$  in spin systems)

submission: use  $\beta$  and  
 $am_q$

M. D'Elia, A. Di Giacomo and C. Pica, Phys. Rev. D **72** (2005) 114510 [arXiv:hep-lat/0503030]

# Conclusions

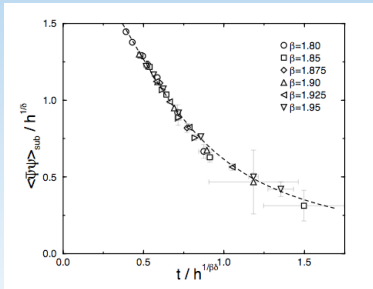
- ▶ Scaling and Finite-Size Scaling functions for  $O(N)$  class
- ▶ relevant FSS region for lattice QCD likely at *smaller* lattices or *smaller*  $m_\pi$
- ▶ large corrections to scaling for large symmetry breaking
- ▶ strategy for comparison

## Outlook

- ▶ comparison to lattice QCD results:  
are calculations in the "interesting" region?
- ▶ comparison to  $O(2)$  scaling functions

# Scaling on the lattice: Staggered vs. Wilson

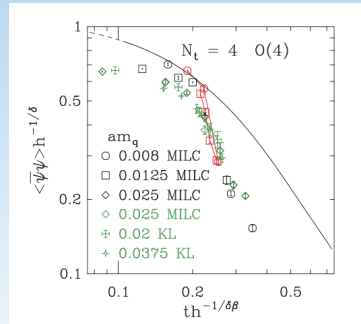
there is a discrepancy in the  $N_f = 2$  scaling behavior between staggered and Wilson fermions:



CP-PACS, A. Ali Khan *et al.*, Phys. Rev. D **63** (2001) 034502

[hep-lat/0008011]

still unresolved!



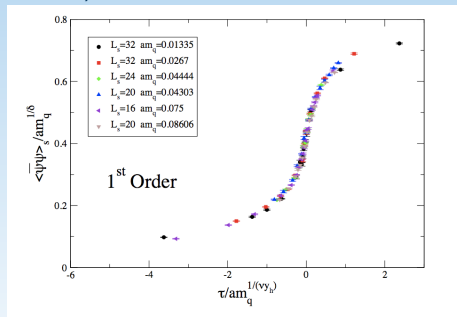
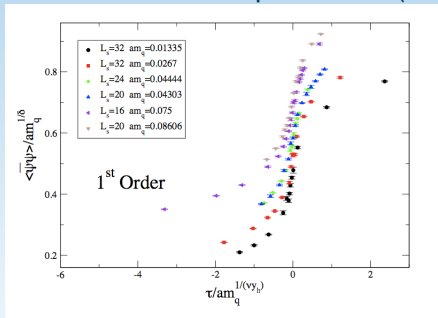
MILC, C. W. Bernard *et al.*, Phys. Rev. D **61** (2000) 054503

[hep-lat/9908008]



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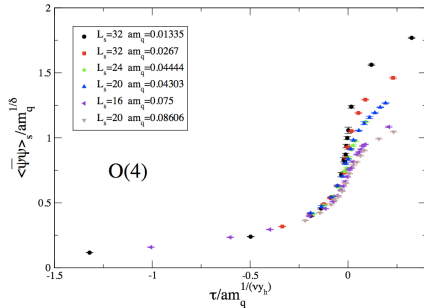
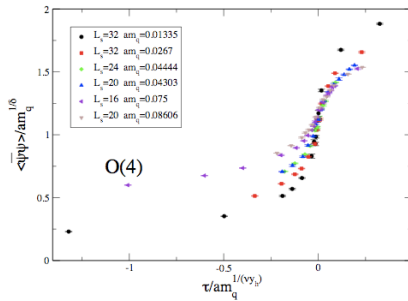
► rescaled order parameter (subtracted) from DiGiacomo et al.



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